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Reducing V_{oc} loss in InGaAsP/InGaAs dual-junction solar cells

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Abstract; Smaller $V_{\alpha \alpha}$ of 1.0 eV/0.75 eV InGaAsP/InGaAs double-junction solar cell(DJSC) than the $V_{\alpha \alpha}$ sum of individual subcells has been observed, and there is little information of the origin of such V_{cc} loss and how to minimize it. In this paper, it is disclosed that the dominant mechanism of minority-carrier transport at back-surfacefield (BSF)/base interface of the bottom subcell is thermionic emission, instead of defect-induced recombination, which is in contrast to previous reports. It also shows that both InP and InAlAs cannot prevent the zinc diffusion effectively. In addition, intermixing of major III-V element occurs as a result of increasing thermal treatment. To suppress the above negative effects, an initial novel InP/InAlAs superlattice (SL) BSF layer is then proposed and employed in bottom InGaAs subcell. The V_{∞} of fabricated cells reach 997.5 mV, and a reduction of 30 mV in Voc loss without lost of Jsc, compared with the results of conventional InP BSF configuration, is achieved. It would benefit the overall V_{∞} for further four-junction solar cells.

Key words: Back-surface field, InGaAsP/InGaAs dual-junction, open-circuit voltage, superlattice. PACS: 88. 40. jp, 78. 67. HC, 75. 40. Mg

InGaAsP/InGaAs双结太阳电池的开路电压损耗抑制

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摘要:现有1.0 eV/0.75 eV InGaAsP/InGaAs双结太阳电池的开路电压小于各子电池的开路电压之和,鲜有研究 探索开路电压损耗的来源以及如何抑制。通过研究发现,InGaAs底电池背场/基区界面处的少数载流子输运 的主要机制是热离子发射,而不是缺陷诱导复合。SIMS测试表明,采用 InP或 InAlAs 背场均不能有效抑制Zn 掺杂剂的扩散。此外,由于生长过程中持续的高温热处理,III-V族主元素在界面处发生了热扩散。为了抑制 上述现象,提出了一种新型InP/InAlAs超晶格背场,并应用到InGaAs底电池中。制备得到的双结太阳电池在 维持短路电流密度不变的情况下,开路电压提升到997.5 mV,与传统采用 InP 背场的双结太阳电池相比,开路 电压损耗降低了30mV。该研究成果对提升四结太阳电池的整体开路电压有重要意义。

关键 词:背场;InGaAsP/InGaAs双结电池;开路电压;超晶格 中图分类号:TM914.4

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Introduction

InGaAsP/InGaAs double-junction solar cells (DJSCs) with approximate bandgap combination of 1.0/ 0.75 eV are used in four-junction configuration to harvest 900~1 700 nm sunlight, and are crucially important for device performances ^[1]. Previous reports Ref. [2-4]

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show that open-circuit voltage (V_{oc}) of InGaAsP/InGaAs DJSC is smaller than the sum of individual subcells. To evaluate solar cells with different bandgaps, bandgap-voltage offset under open-circuit condition (W_{oc}) is introduced ^[5]. For multijunction solar cells, *Woc* can be described as

$$W_{oc}(\mathbf{J}_{sc}) = \frac{1}{q} \left[\sum_{i} E_{g}^{i} - V_{oc} \left(\mathbf{J}_{sc} \right) \right] \qquad , \quad (1)$$

where E_g^i is the bandgap of each subcell, and $V_{\infty}(Jsc)$ stands for the open-circuit voltage when the device produces a given value of short-circuit current (Jsc) under illumination. And the V_{∞} loss for multijunction solar cells would be defined as the gap between W_{∞}^m and the sum of W_{∞}^i of subcells, as $W_{\infty}^m - \sum W_{\infty}^i (J_{sc})$.

 W_{oc}^{m} of InGaAsP/InGaAs DJSC at Jsc of conventional four-junction configuration (about 16.5 mA/cm²) is above 820 mV, higher than the W_{oc} sum of InGaAsP (~ 330 mV) and InGaAs (~340 mV) individual subcells. Part results from previous reports are listed in Table 1, and there is no reference reporting the origin of such V_{oc} loss and how to minimize it.

Experience on III-V semiconductor devices reveals that the diffusion and intermixing at heterojunction interface of InP system always leads to device performance degradation^[6]. Moreover, from the viewpoint of physics of solar cell device, the V_{∞} of solar cells majorly depends on the heterojunction interface between base and backsurface field (BSF) layers. Therefore, considering the thermal history of DJSC structure, the bottom InGaAs subcell, especially the BSF/base interface, might be the key role to reduce V_{∞} loss.

In this paper, the evolution of dopant diffusion and recombination at BSF/base interface with increasing thermal treatment is studied. Based on experimental results, we propose a novel InP/InAlAs superlattice (SL) BSF layer for bottom InGaAs subcell. A reduction of 30 mV in V_{∞} loss is achieved, compared with the results of conventional InP BSF configuration. It shows that such SL BSF would benefit the V_{∞} enhancement for four-junction solar cells.

1 Experiments

Growth are done on *n*-type <100> InP substrates using MOVPE technique. The primary group III and group V precursors used are trimethylgallium (TMGa), trimethylindium (TMIn), trimethylgallium (TMAl), arsine (AsH₃), and phosphine (PH₃). The dopant precursors used are silane (SiH₄) and diethylzinc (DEZn). V/III ratio of 200~300 and growth temperature of 650 °C are em-

ployed, as described previously ^[7].

Г	p+-barrier	100nm	
DH1-	p ⁻ -InGaAsP	500nm	
	p*-barrier	100nm	
L	p ⁺⁺ -InAlGaAs	30nm	
	InGaAs spacer	400nm	
Γ	p ⁺ -barrier	100nm	
DH2	p ⁻ -InGaAsP	500nm	
	p+-barrier	100nm	
L	p ⁺⁺ -InAlGaAs	30nm	
	InGaAs spacer	400nm	
Γ	p+-barrier	100nm	
DH3	p ⁻ -InGaAsP	500nm	
DIIIS	p+-barrier	100nm	
	p**-InAlGaAs	30nm	
	InP sub	350µm	

Fig. 1 Cross-section of MOVPE stack containing three BSF/In-GaAsP/BSF DHs. 图 1 含 3 对 BSF/InGaAsP/BSF 双异质结结构样品示意图

Three periods of isotype p^+ -barrier $(100 \text{ nm})/p^-$ -In- $GaAsP(500nm)/p^+$ -barrier (100nm)/ p^{++} -In (Al_{0.1}Ga_{0.9}) As (100nm) double heterojunctions (DHs), separated by InGaAs spacer layers, are grown in the same stack of MOVPE layers, as illustrated in Fig. 1. Two types of barriers, InP and InAlAs respectively, are employed. After growth, individual DHs are exposed by a series of selective etches. Diluted HCl solution and H₂SO₄: H₂O₂: H₂O solution are used for InP layers and arsenide layers, respectively. The overall element profiles in DHs are obtained through secondary ion mass spectra (SIMS) measurement, while the minority-carrier recombination process in DHs are evaluated using time-resolved photoluminescence (TRPL) technique. It should be pointed out that, the bandgap of p-InGaAsP in DHs is 0.83 eV, for TRPL measurement convenience.

SIMS measurements are performed using Cs⁺ primary beam with a fixed 5 kV acceleration. The positive ions of the quasi-molecular cluster are collected and detected. TRPL measurements with a temporal resolution of ~ 200 ps are performed at room temperature. An H-10330-75 PMT is used to collect PL signals.

The schematic cross-section of InGaAsP/InGaAs (1.0/0.75 eV) DJSC structure is shown in Fig. 2. The active region of each subcell consists of *n*-on-*p* junction (emitter/base) surrounded by *n*-type InP window layer and *p*-type BSF layer. In $(Al_{0.1}Ga_{0.9})$ As tunnel junction is used to connect subcells. The structures are then processed following the standard III-V solar cell device art.

Table 1 Previous reported results for InGaAsP/InGaAs DJSC 表1 文献报道的InGaAsP/InGaAs双结电池电性能

Reference	Method	Bandgap	Illumination	$J_{ m sc}$	$V_{\rm oc}$	$W_{ m oc}$
Reference		Danugap		(mA/cm^2)	(mV)	(mV)
Oshima [2]	MBE	1.0/0.71	AM1.5G	13.1	570	1140
Wu ^[3]	MBE	1.05/0.73	AM1. 5G	16.1	830	950
Zhao [4]	MOVPE	1.07/0.74	AM1. 5D	10. 2	977	833

The cells are 1. 0×1.0 cm² in size.



Fig. 2 Cross-section of the InGaAsP/InGaAs double-junction solar cell structure.

图2 InGaAsP/InGaAs双结太阳电池结构示意图

In-house photovoltaic current density-voltage (J-V) measurements are performed under AMO solar simulator, without GaAs filter. External quantum efficiency (EQE) measurements are performed to give qualitative insight into the spectral response. Cells are placed on 25 °C cooled stages during measurements.

2 Results and discussions

Figure 3 shows the element profiles and PL decay curves for topmost DHs (DH1) in InP-barrier and InAlAs-barrier stacks. The zinc concentration around both DH center regions are of the same level about $5-6 \times 10^{16}$ cm⁻³, which is similar to the typical doping level of base in solar cells. Sharp zinc diffusion profile near the interface between InGaAsP and barriers is observed. It should be pointed out that, the zinc doping level of both InP and InAlAs layers in DHs have been increased to 1- 2×10^{18} cm⁻³, almost one order of magnitude higher than typical doping level used in BSF layer, to exacerbate such diffusion behavior. Moreover, the zinc concentration in the underneath $In(Al_{0.1}Ga_{0.9})$ As layer reaches 2× 10¹⁹ cm⁻³. The nearly identical profiles of zinc atom suggest that both InP and InAlAs layer are of the similar effect as anti-diffusion barriers during the growth.

It evidences in Fig. 3 that the minority-carrier decay time in InAlAs-barrier DH is much larger than that in InP-barrier DH. In symmetrical DHs, the effective lifetime τ_{eff} extracted from PL decay is related to both bulk carrier lifetime τ_{bulk} and surface recombination velocity S, as Ref. [8]

$$\frac{1}{\tau_{\rm eff}} = \frac{1}{\tau_{\rm bulk}} + \frac{2S}{d} \qquad , \quad (2)$$

where d is the thickness of confined layer in the DH. For high-quality materials, the $\tau_{\rm bulk}$ approximately equals to



Fig. 3 Element profiles (a) and PL decay curves (b) for InPbarrier and InAlAs-barrier DH1s.

图 3 InP 和 InAlAs 背场双异质结的元素深度剖析(a)和荧光 发光衰减曲线(b)

the reciprocal product of spontaneous radiative recombination coefficient B and doping concentration ${\cal N}$

$$\tau_{\text{bulk}} = (BN)^{-1} \qquad (3)$$

The value of B could be calculated according to Ref. [9]. With the doping concentration acquired from SIMS, the $\tau_{\rm bulk}$ in DH1 is about 200 ns. By mono-exponential fitting of decay curves, $\tau_{\rm eff}$ of 70 ns and 110 ns are obtained for InP-barrier DH and InAlAs-barrier DH. According to Eq. 2, the experimental S for InP/InGaAsP interface and InAlAs/InGaAsP interface are 232 cm/s and 103 cm/s, respectively. Such small recombination velocity suggests that the radiative process in the bulk dominates the carrier recombination, and the recombination at the interface is nearly neglectable.

The primary mechanism of carrier transport at heterojunction interface includes thermionic emission and defect-induced recombination. When thermionic emission dominates, the recombination current writes ^[10]

$$J_{0} = q \mathrm{S} n_{1} = q \frac{m_{2}}{m_{1}} \left(\frac{2k_{B}T}{\pi m_{1}} \right)^{2} e^{-\frac{\Delta E}{k_{B}T}} \cdot n_{1} \qquad , \quad (4)$$

where m_1 and m_2 are effective mass of confined material and barrier, and ΔE is band offset. It is quite obvious that S would exponentially decreases as the band offset increasing. Notice that electron is the minority-carrier in p-InGaAsP DHs. The values of S for DHs, in the scenario of thermionic emission dominating, are estimated using ΔEc and $m_{1,2}$ from Ref. [11], as listed in Table 2. Surface recombination velocities S of 5793 cm/s and 0.738 cm/s are obtained for InP-barrier DH and InAlAsbarrier DH, respectively. Both calculated and experimental values show that, InAlAs-barrier DH presents smaller surface recombination velocity at heterojunction interface. Consider the complicated carrier transport mechanism at heterojunction interface, the gap between experimental S will not be so large. For example, the sharp diffusion profile of zinc would develop built-in field near the interface, it should reduce the population of minority-carrier reaching the interface and therefore, smaller theoretical value of S could be expected, especially for InP-barrier DH. For InAlAs-barrier DH, the larger experimental S than the calculated S implies the minor existence of trap-induced nonradiative recombination across the interface.

For InGaAs, the conduction band offset ΔEc for InP

Table 2 Calculated surface recombination velocity at barrier/InGaAsP interface using Eq..(3) 表 2 采用公式(3)计算得到的背场/InGaAsP 界面处的表面复 合速率

E	1 达平				
	Barrier	$m_2 \; (m_0)$	$m_1\left(m_0\right)$	$\Delta Ec~(meV)$	S(cm/s)
	InP	0.08	0.047	230	5793
	InAlAs	0.075	0.047	460	0.738

barrier and InAlAs barrier are 0. 25 eV and 0. 52 eV, respectively. The larger offset in conduction band means smaller thermionic emission velocity. Meanwhile, the valence band offset ΔEv for InP barrier and InAlAs barrier are 0. 35 eV and 0. 17 eV^[11]. The smaller offset in valence band indicates lower potential barrier for majoritycarrier. Therefore, InAlAs should be more promising BSF layer in solar cells.

Figure 4 displays the overall SIMS results for DH stacks, and Fig. 5 shows the PL decay curves for individual DHs. Extracted lifetimes are summarized in Table 3. It is obvious that zinc concentration in confined layers rises with increased thermal history from DH1 to DH3. The concentration in DH2 is about 1.0×10^{17} cm⁻³, while the concentration in DH3 is about 2. 0×10^{17} cm⁻³, in spite of the type of barriers. This provides further evidence that both InP and InAlAs present similar ability to block zinc diffusion. Although there is a downward trend from DH1 to DH3, the effective lifetimes in InAlAs-barrier DHs are always longer than those in InP-barrier DHs. It suggests the surface recombination velocity is still dominated by thermionic emission process. With zinc concentration increasing to 1.0×10¹⁷ cm⁻³, the bulk carrier lifetime $\tau_{\rm bulk}$ in DH2 decrease to approximately 100 ns, according to Eq. 2. Therefore, the surface recombination velocity S increase to 434 cm/s and 221 cm/s, for InP-barrier DH2 and InAlAs-barrier DH2 respectively. Since the drop of band offset is the only cause for thermionic emission related increase of S, it is supposed that diffusion of major III-V element across the interface, which would lead to such shrink of band offset, occurs during the thermal treatment.



Fig. 4 Overall SIMS results of (a) InP-barrier DH stack and (b) InAlAs-barrier DH stack. 图 4 InP 双异质结(a)和InAlAs 双异质结(b)的整体 SIMS测试结果

As shown in Table 3, for DHs using the same type of barriers, the $\tau_{\rm eff}$ in DH3 are quite close to the $\tau_{\rm eff}$ in DH2. With zinc concentration of 2. 0×10^{17} cm⁻³, the bulk carrier lifetime $\tau_{\rm bulk}$ in DH3 is approximately 50 ns, and S are 210. 2 cm/s and 55. 6 cm/s, for InP-barrier DH3



Fig. 5 PL decay curves of (a) InP-barrier DHs and (b) InAlAsbarrier DHs.
图 5 InP 双异质结(a)和 InAlAs 双异质结(b)的荧光发光衰减曲线

Table 3 The effective minority-carrier lifetime of the DHs

表3					
	$\tau_{\rm eff}(\rm ns)$	DH1	DH2	DH3	
	InP-barrier	70.0	36.5	35.2	
	InAlAs-barrier	110.0	53.0	45.0	

and InAlAs-barrier DH3 respectively. The abnormal decrease of S is probably due to the photon recycling effect, which leads to the longer $\tau_{\rm eff}$ and smaller S in DH3 than expected.

The steady-state PL of DHs confirms the above hypothesis. As shown in Fig. 6, the PL peak for DH1 and DH2 are of similar intensity, while PL intensities of DH3 are nearly one order of magnitude stronger. Considering the optical configuration of DH stack, the photon recycling is the most effective in DH3, and is suppressed in DH1 and DH2 due to extra absorption from underlying narrow bandgap InGaAs spacers.



Fig. 6 Steady-state PL of (a) InP-barrier DHs and (b) InAlAsbarrier DHs. Weak peaks marked by asteroids in DH1 and DH2 are related to the spacers.

图 6 InP 双异质结(a)和 InAlAs 双异质结(b)的稳态荧光发光曲线。标*的微弱发光峰与背场相关

The results of InGaAsP/InGaAs DJSCs using both InP and InAlAs BSF layers confirm the advantages and effectiveness of InAlAs BSF layer in practical device. Figure 7 shows light J-V and EQE measurements of the devices. Using InAlAs BSF layer, the cell presents an efficiency of 9. 28% with a V_{oc} of 983. 2 mV, a J_{sc} of 15. 6 mA/cm² and an FF of 0. 818. Meanwhile, the device using InP BSF present a V_{oc} of 967. 7 mV, a J_{sc} of 15. 3 mA/



Fig. 7 (a) Light J-V and (b) spectra response curves for In-GaAsP/InGaAs solar cells using InP and InAlAs BSF layers 图 7 采用 InP 和 InAlAs 背场的 InGaAsP/InGaAs 双结电池光 照J-V曲线(a)和量子效率曲线(b)

 $\rm cm^2$ and an FF of 0.819. An enhancement of $V_{\rm oc}$ is obtained, without any cost of J_{s} and FF.

It is well established that SL serve as effective barrier for element diffusion or intermixing, and dislocation threading, and it has been widely used in semiconductor devices such as high electron mobility transistors, laser diodes, electro absorption modulators ^[12-16]. Also, the miniband in SL would not introduce extra potential barrier for carrier transport [17]. An initial five-period InP (2nm)/InAlAs(2nm) SL BSF layer is designed and employed in bottom InGaAs subcell of DJSC. A V_w of 997.5 mV, a Jsc of 15.8 mA/cm² and an FF of 0.824 are obtained as in Fig. 8. Both V_{cc} and Jsc are boosted, as expected, in fabricated SL BSF device. The V_{α} approaches 1.0 V, resulting in a Woc of 752.5 mV. A reduction of 30 mV in V_{α} loss for DJSC is achieved, compared with the conventional InP BSF DJSC.



Fig. 8 Light J-V for InGaAsP/InGaAs DJSC using 5-period InP/ InAlAs SL BSF laver.

图 8 采用 5 对 InP/InAlAs 超晶格背场的 InGaAsP/InGaAs 双结 电池光照 J-V 曲线

3 Conclusions

In general, the use of novel SL BSF layer in the bottom subcell reduces the V_{α} loss in InGaAsP/InGaAs DJSC.

Experiments show that, the mechanism of minoritycarrier transport at BSF/base interface of the bottom subcell of InGaAsP/InGaAs DJSCs is dominated by thermionic emission, instead of defect-induced recombination, which is in contrast to previous reports. It also shows that both InP and InAlAs cannot prevent the zinc diffusion effectively. In addition, intermixing of major III-V element occurs as a result of increasing thermal treatment.

Based on the above results, an initial 5-period InP/ InAlAs SL BSF layer is designed and employed in bottom InGaAs subcell of DJSC. A V. of 997.5 mV, a J. of 15.8 mA/cm² and an FF of 0.824 are obtained. The V_m approaches 1.0 V, resulting in a W_m of 752.5 mV. A reduction of 30 mV in V loss for DJSC is achieved, compared with the results of conventional InP BSF configuration. It suggests that such SL BSF would benefit the V_{ac} enhancement for four-junction solar cells.

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第10届国际应用光学与光子学技术交流大会会议论文征文通知

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Topic1: Advanced Laser Materials and Laser Technology / 新型激光材料与激光器

Topic2: Advanced Laser Processing and Manufacturing / 激光先进制造与装备

Topic3: Laser Transmission and Communication / 激光传输与通信技术

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Topic5: THz Technology and Applications / 太赫兹技术

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Topic21: AI in Optics and Photonics / 人工智能在光学与光子学应用

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